

СЕКЦІЯ «ПРИКЛАДНА МАТЕМАТИКА»

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NON SYMPLIFIED EQUATIONS SYSTEM OF k - ε TURBULENCE MODEL IN CYLINDRICAL COORDINATES

Laminar Navier-Stokes equations in cylindrical coordinates are described in most books on theoretical hydrodynamics [1]-[3]. Reynolds Equations can easy be found in the literature also. But the full equations system of k - ε turbulence model in cylindrical coordinates is very difficult to find. They are usually described in papers only in simplified form. In given paper non simplified system of k - ε turbulence model equations in cylindrical coordinates is derived

Key words: k - ε turbulence model, Reynolds equation, cylindrical coordinates.

1. Reynolds equation in cylindrical coordinates.

Navier-Stokes equations in stresses in cylindrical coordinates have a form

$$\frac{\partial U_r}{\partial t} + U_r \frac{\partial U_r}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_r}{\partial \varphi} + U_z \frac{\partial U_r}{\partial z} - \frac{U_\varphi^2}{r} = -\frac{1}{\rho} \frac{\partial P}{\partial r} - \frac{1}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rr}) + \frac{1}{r} \frac{\partial \tau_{r\varphi}}{\partial \varphi} - \frac{\tau_{\varphi\varphi}}{r} + \frac{\partial \tau_{rz}}{\partial z} \right) \quad (1)$$

$$\frac{\partial U_\varphi}{\partial t} + U_r \frac{\partial U_\varphi}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_\varphi}{\partial \varphi} + U_z \frac{\partial U_\varphi}{\partial z} + \frac{U_\varphi U_r}{r} = -\frac{1}{r} \frac{1}{\rho} \frac{\partial P}{\partial \varphi} - \frac{1}{\rho} \left(\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{r\varphi}) + \frac{1}{r} \frac{\partial \tau_{\varphi\varphi}}{\partial \varphi} + \frac{\partial \tau_{\varphi z}}{\partial z} \right) \quad (2)$$

$$\frac{\partial U_z}{\partial t} + U_r \frac{\partial U_z}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_z}{\partial \varphi} + U_z \frac{\partial U_z}{\partial z} = -\frac{1}{\rho} \frac{\partial P}{\partial z} - \frac{1}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rz}) + \frac{1}{r} \frac{\partial \tau_{\varphi z}}{\partial \varphi} + \frac{\partial \tau_{zz}}{\partial z} \right) \quad (3)$$

$$\frac{1}{r} \frac{\partial r U_r}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial U_z}{\partial z} = \frac{\partial U_r}{\partial r} + \frac{U_r}{r} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial U_z}{\partial z} = 0 \quad (4)$$

Friction stresses τ are given by the formulas

$$\tau_{rr} = -2\mu \frac{\partial U_r}{\partial r}; \quad \tau_{\varphi\varphi} = -2\mu \left(\frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r}{r} \right); \quad \tau_{zz} = -2\mu \frac{\partial U_z}{\partial z}; \quad (5)$$

$$\tau_{r\varphi} = \tau_{\varphi r} = -\mu \left(\frac{1}{r} \frac{\partial U_r}{\partial \varphi} + \frac{\partial U_\varphi}{\partial r} - \frac{U_\varphi}{r} \right); \quad \tau_{rz} = \tau_{zr} = -\mu \left(\frac{\partial U_r}{\partial z} + \frac{\partial U_z}{\partial r} \right); \quad \tau_{\varphi z} = \tau_{z\varphi} = -\mu \left(\frac{1}{r} \frac{\partial U_z}{\partial \varphi} + \frac{\partial U_\varphi}{\partial z} \right); \quad (6)$$

After substituting of (5)-(6) into (1)-(3) right side of (1)-(3) can be transformed to the form

$$\begin{aligned} & \frac{\partial}{\partial r} \frac{1}{r} \mu \frac{\partial r U_r}{\partial r} + \frac{1}{r^2} \frac{\partial}{\partial \varphi} \mu \frac{\partial U_r}{\partial \varphi} - 2\mu \frac{1}{r^2} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial}{\partial z} \mu \frac{\partial U_r}{\partial z} + \frac{\partial}{\partial r} \mu \left(\frac{1}{r} \frac{\partial r U_r}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial U_r}{\partial z} \right) \\ & \left[\frac{\partial}{\partial r} \mu \frac{\partial U_\varphi}{\partial r} + \frac{1}{r} \mu \frac{\partial U_\varphi}{\partial r} - \frac{\mu U_\varphi}{r^2} \right] + \frac{1}{r^2} \frac{\partial}{\partial \varphi} \mu \frac{\partial U_\varphi}{\partial \varphi} + 2 \frac{1}{r^2} \frac{\partial \mu U_r}{\partial \varphi} + \frac{\partial}{\partial z} \mu \frac{\partial U_\varphi}{\partial z} + \frac{1}{r} \frac{\partial}{\partial \varphi} \mu \left(\frac{1}{r} \frac{\partial r U_r}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial U_z}{\partial z} \right) \\ & \left[\frac{\partial}{\partial r} \mu \frac{\partial U_z}{\partial r} + \frac{1}{r} \mu \frac{\partial U_z}{\partial r} \right] + \frac{1}{r^2} \frac{\partial}{\partial \varphi} \mu \frac{\partial U_z}{\partial \varphi} + \frac{\partial}{\partial z} \mu \frac{\partial U_z}{\partial z} + \frac{\partial}{\partial z} \mu \left(\frac{1}{r} \frac{\partial r U_r}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial U_z}{\partial z} \right) \end{aligned}$$

Note, than transformation is correct only if $\mu = const$.

In round brackets the continuity equation is placed, so these brackets can be omitted. The terms in the square brackets are rewritten as single term.

$$\begin{aligned} & \frac{\partial U_r}{\partial t} + U_r \frac{\partial U_r}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_r}{\partial \varphi} + U_z \frac{\partial U_r}{\partial z} - \frac{U_\varphi^2}{r} = \\ & = -\frac{1}{\rho} \frac{\partial P}{\partial r} + \frac{1}{\rho} \left(\frac{\partial}{\partial r} \frac{1}{r} \mu \frac{\partial r U_r}{\partial r} + \frac{1}{r^2} \frac{\partial}{\partial \varphi} \mu \frac{\partial U_r}{\partial \varphi} - 2\mu \frac{1}{r^2} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial}{\partial z} \mu \frac{\partial U_r}{\partial z} \right) \end{aligned} \quad (7)$$

$$\begin{aligned} & \frac{\partial U_\varphi}{\partial t} + U_r \frac{\partial U_\varphi}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_\varphi}{\partial \varphi} + U_z \frac{\partial U_\varphi}{\partial z} + \frac{U_r U_\varphi}{r} = \\ & = -\frac{1}{\rho} \frac{\partial P}{\partial \varphi} + \frac{1}{\rho} \left(\frac{\partial}{\partial r} \frac{1}{r} \mu \frac{\partial r U_\varphi}{\partial r} + \frac{1}{r^2} \frac{\partial}{\partial \varphi} \mu \frac{\partial U_\varphi}{\partial \varphi} + 2 \frac{1}{r^2} \mu \frac{\partial U_r}{\partial \varphi} + \frac{\partial}{\partial z} \mu \frac{\partial U_\varphi}{\partial z} \right) \end{aligned} \quad (8)$$

$$\begin{aligned} & \frac{\partial U_z}{\partial t} + U_r \frac{\partial U_z}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_z}{\partial \varphi} + U_z \frac{\partial U_z}{\partial z} = \\ & = -\frac{1}{\rho} \frac{\partial P}{\partial z} + \frac{1}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} \mu r \frac{\partial U_z}{\partial r} + \frac{\partial}{\partial z} \mu \frac{\partial U_z}{\partial z} + \frac{1}{r^2} \frac{\partial}{\partial \varphi} \mu \frac{\partial U_z}{\partial \varphi} \right) \end{aligned} \quad (9)$$

Now consider as velocity fluctuations forces on flow. For this we rewrite left side of (7)-(9) using continuity equation in conservative form.

$$\begin{aligned} & U_r \frac{\partial U_r}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_r}{\partial \varphi} + U_z \frac{\partial U_r}{\partial z} - \frac{U_\varphi^2}{r} + U_r \left(\frac{\partial U_r}{\partial r} + \frac{\partial U_z}{\partial z} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r}{r} \right) = \\ & = \frac{1}{r} \frac{\partial r U_r U_r}{\partial r} + \frac{1}{r} \frac{\partial U_r U_\varphi}{\partial \varphi} + \frac{\partial U_r U_z}{\partial z} - \frac{U_\varphi^2}{r} \end{aligned} \quad (10)$$

$$\begin{aligned} & U_r \frac{\partial U_\varphi}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r U_\varphi}{r} + U_z \frac{\partial U_\varphi}{\partial z} + U_\varphi \left(\frac{\partial U_r}{\partial r} + \frac{\partial U_z}{\partial z} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r}{r} \right) = \\ & = \frac{1}{r^2} \frac{\partial r^2 U_r U_\varphi}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi U_\varphi}{\partial \varphi} + \frac{\partial U_\varphi U_z}{\partial z} \end{aligned} \quad (11)$$

$$\begin{aligned}
& U_r \frac{\partial U_z}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_z}{\partial \varphi} + U_z \frac{\partial U_z}{\partial z} + U_z \left(\frac{\partial U_r}{\partial r} + \frac{\partial U_z}{\partial z} + \frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r}{r} \right) = \\
& = \frac{1}{r} \frac{\partial r U_r U_z}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi U_z}{\partial \varphi} + \frac{\partial U_z U_z}{\partial z}
\end{aligned} \tag{12}$$

Now we decomposed instantaneous quantities into sum of time-averaged and fluctuating parts, substitute this sum into (10)-(12) and average result. After this expressions for averaged velocities rewrite in original, non conservative form.

$$\begin{aligned}
& U_r \frac{\partial U_z}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_z}{\partial \varphi} + U_z \frac{\partial U_z}{\partial z} = \frac{1}{r} \frac{\partial r U_r U_z}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi U_z}{\partial \varphi} + \frac{\partial U_z U_z}{\partial z} = \\
& = \overline{U_r} \frac{\partial \overline{U_z}}{\partial r} + \frac{\overline{U_\varphi}}{r} \frac{\partial \overline{U_z}}{\partial \varphi} + \overline{U_z} \frac{\partial \overline{U_z}}{\partial z} + \frac{1}{r} \frac{\partial \overline{r u_r u_z}}{\partial r} + \frac{1}{r} \frac{\partial \overline{u_\varphi u_z}}{\partial \varphi} + \frac{\partial \overline{u_z u_z}}{\partial z}
\end{aligned} \tag{13}$$

$$\begin{aligned}
& U_r \frac{\partial U_\varphi}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r U_\varphi}{r} + U_z \frac{\partial U_\varphi}{\partial z} = \frac{1}{r^2} \frac{\partial r^2 U_r U_\varphi}{\partial r} + \frac{1}{r} \frac{\partial U_\varphi U_\varphi}{\partial \varphi} + \frac{\partial U_\varphi U_z}{\partial z} = \\
& = \overline{U_r} \frac{\partial \overline{U_\varphi}}{\partial r} + \frac{\overline{U_\varphi}}{r} \frac{\partial \overline{U_\varphi}}{\partial \varphi} + \frac{\overline{U_r U_\varphi}}{r} + \overline{U_z} \frac{\partial \overline{U_\varphi}}{\partial z} + \frac{1}{r^2} \frac{\partial r^2 \overline{u_r u_\varphi}}{\partial r} + \frac{1}{r} \frac{\partial \overline{u_\varphi u_\varphi}}{\partial \varphi} + \frac{\partial \overline{u_\varphi u_z}}{\partial z}
\end{aligned} \tag{14}$$

$$\begin{aligned}
& U_r \frac{\partial U_r}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_r}{\partial \varphi} + U_z \frac{\partial U_r}{\partial z} - \frac{U_\varphi^2}{r} = \frac{1}{r} \frac{\partial r U_r U_r}{\partial r} + \frac{1}{r} \frac{\partial U_r U_\varphi}{\partial \varphi} + \frac{\partial U_r U_z}{\partial z} - \frac{U_\varphi^2}{r} = \\
& = \frac{\overline{U_\varphi}}{r} \frac{\partial \overline{U_r}}{\partial \varphi} + \overline{U_r} \frac{\partial \overline{U_r}}{\partial r} + \overline{U_z} \frac{\partial \overline{U_r}}{\partial z} - \frac{\overline{U_\varphi^2}}{r} + \frac{1}{r} \frac{\partial \overline{r u_r u_r}}{\partial r} + \frac{1}{r} \frac{\partial \overline{u_\varphi u_\varphi}}{\partial \varphi} + \frac{\partial \overline{u_r u_z}}{\partial z} - \frac{\overline{u_\varphi u_\varphi}}{r}
\end{aligned} \tag{15}$$

Now substitute (13)-(15) into (1)-(3). After simple transformation we obtain:

$$\begin{aligned}
& \frac{\partial \overline{U_r}}{\partial t} + \overline{U_r} \frac{\partial \overline{U_r}}{\partial r} + \frac{\overline{U_\varphi}}{r} \frac{\partial \overline{U_r}}{\partial \varphi} + \overline{U_z} \frac{\partial \overline{U_r}}{\partial z} - \frac{\overline{U_\varphi^2}}{r} = \\
& = -\frac{1}{\rho} \frac{\partial \overline{P}}{\partial r} + \frac{1}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} (r (\tau_{rr} - \overline{u_r u_r})) \right) + \frac{1}{r} \frac{\partial}{\partial \varphi} (\tau_{r\varphi} - \overline{u_r u_\varphi}) - \frac{(\tau_{\varphi\varphi} - \overline{u_\varphi u_\varphi})}{r} + \frac{\partial}{\partial z} (\tau_{rz} - \overline{u_r u_z})
\end{aligned}$$

$$\begin{aligned}
& \frac{\partial U_\varphi}{\partial t} + U_r \frac{\partial U_\varphi}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_\varphi}{\partial \varphi} + U_z \frac{\partial U_\varphi}{\partial z} + \frac{U_\varphi U_r}{r} = \\
& = -\frac{1}{\rho} \frac{\partial P}{\partial \varphi} + \frac{1}{\rho} \left(\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 (\tau_{r\varphi} - \overline{u_r u_\varphi})) \right) + \frac{1}{r} \frac{\partial}{\partial \varphi} (\tau_{\varphi\varphi} - \overline{u_\varphi u_\varphi}) + \frac{\partial}{\partial z} (\tau_{\varphi z} - \overline{u_\varphi u_z})
\end{aligned}$$

$$\begin{aligned}
& \frac{\partial U_z}{\partial t} + U_r \frac{\partial U_z}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_z}{\partial \varphi} + U_z \frac{\partial U_z}{\partial z} = \\
& = -\frac{1}{\rho} \frac{\partial P}{\partial z} + \frac{1}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} (r (\tau_{rz} - \overline{u_r u_z})) \right) + \frac{1}{r} \frac{\partial}{\partial \varphi} (\tau_{\varphi z} - \overline{u_\varphi u_z}) + \frac{\partial}{\partial z} (\tau_{zz} - \overline{u_z u_z})
\end{aligned}$$

Following Boussinesq hypothesis we introduce turbulent viscosity as:

$$\begin{aligned} -\overline{u_r u_r} &= 2\mu \frac{\partial U_r}{\partial r}; \quad -\overline{u_\varphi u_\varphi} = 2\mu_t \left(\frac{1}{r} \frac{\partial U_\varphi}{\partial \varphi} + \frac{U_r}{r} \right); \quad -\overline{u_z u_z} = 2\mu_t \frac{\partial U_z}{\partial z}; \\ -\overline{u_r u_\varphi} &= \mu_t \left(\frac{1}{r} \frac{\partial U_r}{\partial \varphi} + \frac{\partial U_\varphi}{\partial r} - \frac{U_\varphi}{r} \right); \quad -\overline{u_r u_z} = \mu_t \left(\frac{\partial U_r}{\partial z} + \frac{\partial U_z}{\partial r} \right); \quad -\overline{u_\varphi u_z} = \mu_t \left(\frac{1}{r} \frac{\partial U_z}{\partial \varphi} + \frac{\partial U_\varphi}{\partial z} \right). \end{aligned}$$

Then turbulent Navier-Stokes or Reynolds equations can be written in the form:

$$\begin{aligned} \frac{\partial U_r}{\partial t} + U_r \frac{\partial U_r}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_r}{\partial \varphi} + U_z \frac{\partial U_r}{\partial z} - \frac{U_\varphi^2}{r} = \\ = -\frac{1}{\rho} \frac{\partial P}{\partial r} + \frac{1}{\rho} \left(\frac{\partial}{\partial r} \frac{1}{r} (\mu + \mu_t) \frac{\partial r U_r}{\partial r} + \frac{1}{r^2} \frac{\partial}{\partial \varphi} (\mu + \mu_t) \frac{\partial U_r}{\partial \varphi} - 2(\mu + \mu_t) \frac{1}{r^2} \frac{\partial U_\varphi}{\partial \varphi} + \frac{\partial}{\partial z} (\mu + \mu_t) \frac{\partial U_r}{\partial z} \right) \end{aligned} \quad (16)$$

$$\begin{aligned} \frac{\partial U_\varphi}{\partial t} + U_r \frac{\partial U_\varphi}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_\varphi}{\partial \varphi} + U_z \frac{\partial U_\varphi}{\partial z} + \frac{U_r U_\varphi}{r} = \\ = -\frac{1}{\rho} \frac{1}{r} \frac{\partial P}{\partial \varphi} + \frac{1}{\rho} \left(\frac{\partial}{\partial r} \frac{1}{r} (\mu + \mu_t) \frac{\partial r U_\varphi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial}{\partial \varphi} (\mu + \mu_t) \frac{\partial U_\varphi}{\partial \varphi} + 2 \frac{1}{r^2} (\mu + \mu_t) \frac{\partial U_r}{\partial \varphi} + \frac{\partial}{\partial z} (\mu + \mu_t) \frac{\partial U_\varphi}{\partial z} \end{aligned} \quad (17)$$

$$\begin{aligned} \frac{\partial U_z}{\partial t} + U_r \frac{\partial U_z}{\partial r} + \frac{U_\varphi}{r} \frac{\partial U_z}{\partial \varphi} + U_z \frac{\partial U_z}{\partial z} = \\ = -\frac{1}{\rho} \frac{\partial P}{\partial z} + \frac{1}{\rho} \left(\frac{1}{r} \frac{\partial}{\partial r} (\mu + \mu_t) r \frac{\partial U_z}{\partial r} + \frac{\partial}{\partial z} (\mu + \mu_t) \frac{\partial U_z}{\partial z} + \frac{1}{r^2} \frac{\partial}{\partial \varphi} (\mu + \mu_t) \frac{\partial U_z}{\partial \varphi} \right) \end{aligned} \quad (18)$$

2. Turbulent kinetic energy transfer, convective and pressure terms.

Now we decomposed instantaneous velocities and pressure as sum of time-averaged and fluctuating parts $U_r = \overline{U_r} + u_r$, $U_\varphi = \overline{U_\varphi} + u_\varphi$, $U_z = \overline{U_z} + u_z$, $\overline{P} = P + p$, substitute this sum into (1)-(3) and average result. Diffusion terms will be transformed separately.

$$\begin{aligned} u_\varphi \left(\frac{\partial (\overline{U_\varphi} + u_\varphi)}{\partial t} + (\overline{U_r} + u_r) \frac{\partial (\overline{U_\varphi} + u_\varphi)}{\partial r} + \frac{(\overline{U_\varphi} + u_\varphi)}{r} \frac{\partial (\overline{U_\varphi} + u_\varphi)}{\partial \varphi} + (\overline{U_z} + u_z) \frac{\partial (\overline{U_\varphi} + u_\varphi)}{\partial z} + \frac{(\overline{U_r} + u_r)(\overline{U_\varphi} + u_\varphi)}{r} + \right. \\ \left. + \frac{1}{\rho} \frac{1}{r} \frac{\partial (\overline{P} + p)}{\partial \varphi} \right) = \left(\frac{\partial (0.5 \overline{u_\varphi^2})}{\partial t} + \overline{U_r} \frac{\partial (0.5 \overline{u_\varphi^2})}{\partial r} + \frac{\overline{U_\varphi}}{r} \frac{\partial (0.5 \overline{u_\varphi^2})}{\partial \varphi} + \overline{U_z} \frac{\partial (0.5 \overline{u_\varphi^2})}{\partial z} \right) + \left(u_\varphi u_r \frac{\partial \overline{U_\varphi}}{\partial r} + u_\varphi u_\varphi \frac{1}{r} \frac{\partial \overline{U_\varphi}}{\partial \varphi} + \right. \\ \left. + u_\varphi u_z \frac{\partial \overline{U_\varphi}}{\partial z} \right) = \left(u_\varphi u_r \frac{\partial u_\varphi}{\partial r} + u_\varphi u_\varphi \frac{1}{r} \frac{\partial u_\varphi}{\partial \varphi} + u_\varphi u_z \frac{\partial u_\varphi}{\partial z} \right) + \left(u_\varphi u_\varphi \frac{\overline{U_r}}{r} + u_\varphi u_r \frac{\overline{U_\varphi}}{r} + \frac{u_r u_\varphi u_\varphi}{r} \right) + \frac{1}{\rho} \frac{1}{r} u_\varphi \frac{\partial p}{\partial \varphi} \end{aligned} \quad (19)$$

$$\begin{aligned}
 & u_r \left(\frac{\partial(\overline{U_r + u_r})}{\partial t} + (\overline{U_r + u_r}) \frac{\partial(\overline{U_r + u_r})}{\partial r} + \frac{(\overline{U_\phi + u_\phi})}{r} \frac{\partial(\overline{U_r + u_r})}{\partial \phi} + (\overline{U_z + u_z}) \frac{\partial(\overline{U_r + u_r})}{\partial z} - \frac{(\overline{U_\phi + u_\phi})^2}{r} + \frac{1}{\rho} \frac{\partial(\overline{P} + p)}{\partial r} \right) = \\
 & = \left(\frac{\partial(0.5u_r^2)}{\partial t} \right) + \left(\overline{U_r} \frac{\partial(0.5u_r^2)}{\partial r} + u_r u_r \frac{\partial \overline{U_r}}{\partial r} + u_r u_r \frac{\partial u_r}{\partial r} \right) + \left(\frac{\overline{U_\phi}}{r} \frac{\partial(0.5u_r^2)}{\partial \phi} + u_r u_\phi \frac{1}{r} \frac{\partial \overline{U_r}}{\partial \phi} + u_r u_\phi \frac{1}{r} \frac{\partial u_r}{\partial \phi} \right) + \\
 & + \left(\overline{U_z} \frac{\partial(0.5u_r^2)}{\partial z} + u_r u_z \frac{\partial \overline{U_r}}{\partial z} + u_r u_z \frac{\partial u_r}{\partial z} \right) - \left(u_r u_\phi \frac{\overline{U_\phi}}{r} + u_r u_\phi \frac{\overline{U_\phi}}{r} + \frac{u_r u_\phi u_\phi}{r} \right) + \frac{1}{\rho} u_r \frac{\partial p}{\partial r}
 \end{aligned} \tag{20}$$

$$\begin{aligned}
 & u_z \left(\frac{\partial(\overline{U_z + u_z})}{\partial t} + (\overline{U_r + u_r}) \frac{\partial(\overline{U_z + u_z})}{\partial r} + \frac{(\overline{U_\phi + u_\phi})}{r} \frac{\partial(\overline{U_z + u_z})}{\partial \phi} + (\overline{U_z + u_z}) \frac{\partial(\overline{U_z + u_z})}{\partial z} + \frac{1}{\rho} \frac{\partial(\overline{P} + p)}{\partial z} \right) = \\
 & = \left(\frac{\partial(0.5u_z^2)}{\partial t} \right) + \left(\overline{U_r} \frac{\partial(0.5u_z^2)}{\partial r} + u_z u_r \frac{\partial \overline{U_z}}{\partial r} + u_z u_r \frac{\partial u_z}{\partial r} \right) + \left(\frac{\overline{U_\phi}}{r} \frac{\partial(0.5u_z^2)}{\partial \phi} + u_z u_\phi \frac{1}{r} \frac{\partial \overline{U_z}}{\partial \phi} + u_z u_\phi \frac{1}{r} \frac{\partial u_z}{\partial \phi} \right) + \\
 & + \left(\overline{U_z} \frac{\partial(0.5u_z^2)}{\partial z} + u_z u_z \frac{1}{r} \frac{\partial \overline{U_z}}{\partial z} + u_z u_z \frac{1}{r} \frac{\partial u_z}{\partial z} \right) + \frac{1}{\rho} u_z \frac{\partial p}{\partial z}
 \end{aligned} \tag{21}$$

Summing (19)-(21) we obtain part of equation of turbulent kinetic energy transfer.

$$\begin{aligned}
 & \frac{\partial k}{\partial t} + \overline{U_r} \frac{\partial k}{\partial r} + \frac{\overline{U_\phi}}{r} \frac{\partial k}{\partial \phi} + \overline{U_z} \frac{\partial k}{\partial z} + \left[\overline{u_r u_r} \frac{\partial \overline{U_r}}{\partial r} + \overline{u_\phi u_\phi} \left(\frac{1}{r} \frac{\partial \overline{U_\phi}}{\partial \phi} + \frac{\overline{U_r}}{r} \right) + \overline{u_z u_z} \frac{\partial \overline{U_z}}{\partial z} + \overline{u_\phi u_r} \left(\frac{\partial \overline{U_\phi}}{\partial r} + \frac{1}{r} \frac{\partial \overline{U_r}}{\partial \phi} - \frac{\overline{U_\phi}}{r} \right) \right. \\
 & + \overline{u_z u_r} \left(\frac{\partial \overline{U_z}}{\partial r} + \frac{\partial \overline{U_r}}{\partial z} \right) + \overline{u_\phi u_z} \left(\frac{1}{r} \frac{\partial \overline{U_z}}{\partial \phi} + \frac{\partial \overline{U_\phi}}{\partial z} \right) \left. \right] + \overline{u_r u_r} \frac{\partial u_r}{\partial r} + \overline{u_\phi u_\phi} \frac{1}{r} \frac{\partial u_\phi}{\partial \phi} + \overline{u_z u_z} \frac{\partial u_z}{\partial z} + \overline{u_\phi u_r} \left(\frac{\partial u_\phi}{\partial r} + \frac{1}{r} \frac{\partial u_z}{\partial \phi} \right) + \\
 & + \overline{u_z u_r} \left(\frac{\partial u_z}{\partial r} + \frac{\partial u_r}{\partial z} \right) + \overline{u_\phi u_z} \left(\frac{1}{r} \frac{\partial u_\phi}{\partial \phi} + \frac{\partial u_\phi}{\partial z} \right) - \frac{1}{\rho} \left(\frac{1}{r} u_\phi \frac{\partial p}{\partial \phi} + u_r \frac{\partial p}{\partial r} + u_z \frac{\partial p}{\partial z} \right)
 \end{aligned}$$

The sum can be written in the form (22). Sense of terms will be described later.

$$\frac{\partial k}{\partial t} + \overline{U_r} \frac{\partial k}{\partial r} + \frac{\overline{U_\phi}}{r} \frac{\partial k}{\partial \phi} + \overline{U_z} \frac{\partial k}{\partial z} - P - Diffk_{urb} \tag{22}$$

$$\begin{aligned}
 P & = -\overline{u_r u_r} \frac{\partial \overline{U_r}}{\partial r} - \overline{u_\phi u_\phi} \left(\frac{1}{r} \frac{\partial \overline{U_\phi}}{\partial \phi} + \frac{\overline{U_r}}{r} \right) - \overline{u_z u_z} \frac{\partial \overline{U_z}}{\partial z} - \overline{u_\phi u_r} \left(\frac{\partial \overline{U_\phi}}{\partial r} + \frac{1}{r} \frac{\partial \overline{U_r}}{\partial \phi} - \frac{\overline{U_\phi}}{r} \right) - \\
 & - \overline{u_z u_r} \left(\frac{\partial \overline{U_z}}{\partial r} + \frac{\partial \overline{U_r}}{\partial z} \right) - \overline{u_\phi u_z} \left(\frac{1}{r} \frac{\partial \overline{U_z}}{\partial \phi} + \frac{\partial \overline{U_\phi}}{\partial z} \right) = \\
 & = \nu_t \left[2 \left(\frac{\partial \overline{U_r}}{\partial r} \right)^2 + 2 \left(\frac{1}{r} \frac{\partial \overline{U_\phi}}{\partial \phi} + \frac{\overline{U_r}}{r} \right)^2 + 2 \left(\frac{\partial \overline{U_z}}{\partial z} \right)^2 + \left(\frac{\partial \overline{U_\phi}}{\partial r} + \frac{1}{r} \frac{\partial \overline{U_r}}{\partial \phi} - \frac{\overline{U_\phi}}{r} \right)^2 + \left(\frac{\partial \overline{U_z}}{\partial r} + \frac{\partial \overline{U_r}}{\partial z} \right)^2 + \left(\frac{1}{r} \frac{\partial \overline{U_z}}{\partial \phi} + \frac{\partial \overline{U_\phi}}{\partial z} \right)^2 \right]
 \end{aligned} \tag{23}$$

$$\begin{aligned}
 \overline{Diffk_{turb}} = & -\overline{u_r u_r \frac{\partial u_r}{\partial r}} - \overline{u_\varphi u_\varphi \frac{1}{r} \frac{\partial u_\varphi}{\partial \varphi}} - \overline{u_z u_z \frac{\partial u_z}{\partial z}} - \overline{u_\varphi u_r \left(\frac{\partial u_\varphi}{\partial r} + \frac{1}{r} \frac{\partial u_z}{\partial \varphi} \right)} - \\
 & - \overline{u_z u_r \left(\frac{\partial u_z}{\partial r} + \frac{\partial u_r}{\partial z} \right)} - \overline{u_\varphi u_z \left(\frac{1}{r} \frac{\partial u_\varphi}{\partial \varphi} + \frac{\partial u_\varphi}{\partial z} \right)} - \frac{1}{\rho} \left(\overline{\frac{1}{r} u_\varphi \frac{\partial p}{\partial \varphi}} + \overline{u_r \frac{\partial p}{\partial r}} + \overline{u_z \frac{\partial p}{\partial z}} \right)
 \end{aligned} \tag{24}$$

3. Turbulent kinetic energy transfer, viscous diffusion terms.

Now we transform diffusion terms of (1)-(3).

$$\begin{aligned}
 & \frac{\mu}{\rho} u_r \left(\frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r (\overline{U_r} + u_r)}{\partial r} + \frac{1}{r^2} \frac{\partial^2 (\overline{U_r} + u_r)}{\partial \varphi^2} - 2 \frac{1}{r^2} \frac{\partial (\overline{U_\varphi} + u_\varphi)}{\partial \varphi} + \frac{\partial^2 (\overline{U_r} + u_r)}{\partial z^2} \right) = \\
 & = \frac{\mu}{\rho} \left(u_r \frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r u_r}{\partial r} + \frac{1}{r^2} u_r \frac{\partial^2 u_r}{\partial \varphi^2} - 2 \frac{1}{r^2} u_r \frac{\partial u_\varphi}{\partial \varphi} + u_r \frac{\partial^2 u_r}{\partial z^2} \right) \\
 & \frac{\mu}{\rho} u_\varphi \left(\frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r (\overline{U_\varphi} + u_\varphi)}{\partial r} + \frac{1}{r^2} \frac{\partial^2 (\overline{U_\varphi} + u_\varphi)}{\partial \varphi^2} + 2 \frac{1}{r^2} \frac{\partial (\overline{U_r} + u_r)}{\partial \varphi} + \frac{\partial^2 (\overline{U_\varphi} + u_\varphi)}{\partial z^2} \right) = \\
 & = \frac{\mu}{\rho} \left(u_\varphi \frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r u_\varphi}{\partial r} + \frac{1}{r^2} u_\varphi \frac{\partial^2 u_\varphi}{\partial \varphi^2} + 2 \frac{1}{r^2} u_\varphi \frac{\partial u_r}{\partial \varphi} + u_\varphi \frac{\partial^2 u_\varphi}{\partial z^2} \right) \\
 & \frac{\mu}{\rho} u_z \left(\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial (\overline{U_z} + u_z)}{\partial r} + \frac{1}{r^2} \frac{\partial^2 (\overline{U_z} + u_z)}{\partial \varphi^2} + \frac{\partial^2 (\overline{U_z} + u_z)}{\partial z^2} \right) = \\
 & = \frac{\mu}{\rho} \left(u_z \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial u_z}{\partial r} + \frac{1}{r^2} u_z \frac{\partial^2 u_z}{\partial \varphi^2} + u_z \frac{\partial^2 u_z}{\partial z^2} \right)
 \end{aligned}$$

Note, that after opening the brackets derivatives of averaged velocities are found to be multiplied on fluctuation terms. As a result of averaging such products are equal to zero.

Using following formulas:

$$u \frac{\partial^2 u}{\partial \xi^2} = 0.5 \frac{\partial u^2}{\partial \xi^2} - \frac{\partial u}{\partial \xi} \frac{\partial u}{\partial \xi};$$

$$u \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial u}{\partial r} = 0.5 \frac{\partial^2 u^2}{\partial r^2} - \frac{\partial u}{\partial r} \frac{\partial u}{\partial r} + \frac{1}{r} 0.5 \frac{\partial u^2}{\partial r};$$

$$u \frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r u}{\partial r} = 0.5 \frac{\partial^2 u^2}{\partial r^2} - \frac{\partial u}{\partial r} \frac{\partial u}{\partial r} + \frac{1}{r} 0.5 \frac{\partial u^2}{\partial r} - \frac{u^2}{r^2}$$

we obtain:

$$\begin{aligned}
 & u_r \frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r u_r}{\partial r} + \frac{1}{r^2} u_r \frac{\partial^2 u_r}{\partial \varphi^2} - 2 \frac{1}{r^2} u_r \frac{\partial u_\varphi}{\partial \varphi} + u_r \frac{\partial^2 u_r}{\partial z^2} = \\
 & = 0.5 \frac{\partial^2 u_r^2}{\partial r^2} - \frac{\partial u_r}{\partial r} \frac{\partial u_r}{\partial r} + \frac{1}{r} 0.5 \frac{\partial u_r^2}{\partial r} - \frac{u_r^2}{r^2} + \frac{1}{r^2} 0.5 \frac{\partial u_r^2}{\partial \varphi^2} - \frac{1}{r^2} \frac{\partial u_r}{\partial \varphi} \frac{\partial u_r}{\partial \varphi} - 2 \frac{1}{r^2} u_r \frac{\partial u_\varphi}{\partial \varphi} + 0.5 \frac{\partial^2 u_r^2}{\partial z^2} - \frac{\partial u_r}{\partial z} \frac{\partial u_r}{\partial z}
 \end{aligned}$$

$$\begin{aligned}
& u_\varphi \frac{\partial}{\partial r} \frac{1}{r} \frac{\partial r u_\varphi}{\partial r} + \frac{1}{r^2} u_\varphi \frac{\partial^2 u_\varphi}{\partial \varphi^2} + 2 \frac{1}{r^2} u_\varphi \frac{\partial u_r}{\partial \varphi} + u_\varphi \frac{\partial^2 u_\varphi}{\partial z^2} = \\
& = 0.5 \frac{\partial^2 u_\varphi^2}{\partial r^2} - \frac{\partial u_\varphi}{\partial r} \frac{\partial u_\varphi}{\partial r} + \frac{1}{r} 0.5 \frac{\partial u_\varphi^2}{\partial r} - \frac{u_\varphi^2}{r^2} + \frac{1}{r^2} 0.5 \frac{\partial u_\varphi^2}{\partial \varphi^2} - \frac{1}{r^2} \frac{\partial u_\varphi}{\partial \varphi} \frac{\partial u_\varphi}{\partial \varphi} + 2 \frac{1}{r^2} u_\varphi \frac{\partial u_r}{\partial \varphi} + 0.5 \frac{\partial^2 u_\varphi^2}{\partial z^2} - \frac{\partial u_\varphi}{\partial z} \frac{\partial u_\varphi}{\partial z} \\
& u_z \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial u_z}{\partial r} + \frac{1}{r^2} u_z \frac{\partial^2 u_z}{\partial \varphi^2} + u_z \frac{\partial^2 u_z}{\partial z^2} = \\
& = 0.5 \frac{\partial^2 u_z^2}{\partial r^2} - \frac{\partial u_z}{\partial r} \frac{\partial u_z}{\partial r} + \frac{1}{r} 0.5 \frac{\partial u_z^2}{\partial r} + \frac{1}{r^2} 0.5 \frac{\partial u_z^2}{\partial \varphi^2} - \frac{1}{r^2} \frac{\partial u_z}{\partial \varphi} \frac{\partial u_z}{\partial \varphi} + 0.5 \frac{\partial^2 u_z^2}{\partial z^2} - \frac{\partial u_z}{\partial z} \frac{\partial u_z}{\partial z}.
\end{aligned}$$

Summing and regrouping terms we obtain:

$$\begin{aligned}
& 0.5 \frac{\partial^2 u_r^2}{\partial r^2} + \frac{1}{r} 0.5 \frac{\partial u_r^2}{\partial r} + \frac{1}{r^2} 0.5 \frac{\partial u_r^2}{\partial \varphi^2} + 0.5 \frac{\partial^2 u_r^2}{\partial z^2} - \frac{\partial u_r}{\partial r} \frac{\partial u_r}{\partial r} - \frac{1}{r^2} \frac{\partial u_r}{\partial \varphi} \frac{\partial u_r}{\partial \varphi} - \frac{\partial u_r}{\partial z} \frac{\partial u_r}{\partial z} - 2 \frac{1}{r^2} u_r \frac{\partial u_\varphi}{\partial \varphi} - \frac{u_r^2}{r^2} + \\
& + 0.5 \frac{\partial^2 u_\varphi^2}{\partial r^2} + \frac{1}{r} 0.5 \frac{\partial u_\varphi^2}{\partial r} + \frac{1}{r^2} 0.5 \frac{\partial u_\varphi^2}{\partial \varphi^2} + 0.5 \frac{\partial^2 u_\varphi^2}{\partial z^2} - \frac{\partial u_\varphi}{\partial r} \frac{\partial u_\varphi}{\partial r} - \frac{1}{r^2} \frac{\partial u_\varphi}{\partial \varphi} \frac{\partial u_\varphi}{\partial \varphi} - \frac{\partial u_\varphi}{\partial z} \frac{\partial u_\varphi}{\partial z} + 2 \frac{1}{r^2} u_\varphi \frac{\partial u_r}{\partial \varphi} - \frac{u_\varphi^2}{r^2} + \\
& + 0.5 \frac{\partial^2 u_z^2}{\partial r^2} + \frac{1}{r} 0.5 \frac{\partial u_z^2}{\partial r} + \frac{1}{r^2} 0.5 \frac{\partial u_z^2}{\partial \varphi^2} + 0.5 \frac{\partial^2 u_z^2}{\partial z^2} - \frac{\partial u_z}{\partial r} \frac{\partial u_z}{\partial r} - \frac{1}{r^2} \frac{\partial u_z}{\partial \varphi} \frac{\partial u_z}{\partial \varphi} - \frac{\partial u_z}{\partial z} \frac{\partial u_z}{\partial z} = \\
& = \frac{\partial^2 k}{\partial r^2} + \frac{1}{r} \frac{\partial k}{\partial r} + \frac{1}{r^2} \frac{\partial k}{\partial \varphi^2} + \frac{\partial^2 k}{\partial z^2} - \varepsilon = \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial k}{\partial r} + \frac{1}{r^2} \frac{\partial k}{\partial \varphi^2} + \frac{\partial^2 k}{\partial z^2} - \varepsilon
\end{aligned}$$

Here:

$$\begin{aligned}
\varepsilon = & - \frac{\partial u_r}{\partial r} \frac{\partial u_r}{\partial r} - \frac{1}{r^2} \frac{\partial u_r}{\partial \varphi} \frac{\partial u_r}{\partial \varphi} - \frac{\partial u_r}{\partial z} \frac{\partial u_r}{\partial z} - 2 \frac{1}{r^2} u_r \frac{\partial u_\varphi}{\partial \varphi} - \frac{u_r^2}{r^2} - \frac{\partial u_\varphi}{\partial r} \frac{\partial u_\varphi}{\partial r} - \frac{1}{r^2} \frac{\partial u_\varphi}{\partial \varphi} \frac{\partial u_\varphi}{\partial \varphi} - \frac{\partial u_\varphi}{\partial z} \frac{\partial u_\varphi}{\partial z} + \\
& + 2 \frac{1}{r^2} u_\varphi \frac{\partial u_r}{\partial \varphi} - \frac{u_\varphi^2}{r^2} - \frac{\partial u_z}{\partial r} \frac{\partial u_z}{\partial r} - \frac{1}{r^2} \frac{\partial u_z}{\partial \varphi} \frac{\partial u_z}{\partial \varphi} - \frac{\partial u_z}{\partial z} \frac{\partial u_z}{\partial z}
\end{aligned}$$

Collecting all terms of all equation we obtain:

$$\begin{aligned}
\frac{\partial k}{\partial t} + \overline{U_r} \frac{\partial k}{\partial r} + \frac{\overline{U_\varphi}}{r} \frac{\partial k}{\partial \varphi} + \overline{U_z} \frac{\partial k}{\partial z} & = \nu \left(\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial k}{\partial r} + \frac{1}{r^2} \frac{\partial k}{\partial \varphi^2} + \frac{\partial^2 k}{\partial z^2} \right) + Diffk_{urb} + P - \varepsilon \\
\frac{\partial \varepsilon}{\partial t} + \overline{U_r} \frac{\partial \varepsilon}{\partial r} + \frac{\overline{U_\varphi}}{r} \frac{\partial \varepsilon}{\partial \varphi} + \overline{U_z} \frac{\partial \varepsilon}{\partial z} & = \nu \left(\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial \varepsilon}{\partial r} + \frac{1}{r^2} \frac{\partial \varepsilon}{\partial \varphi^2} + \frac{\partial^2 \varepsilon}{\partial z^2} \right) + Diff\varepsilon_{urb} + \frac{\varepsilon}{k} (C_1 P - C_2 \varepsilon) \\
\nu_t & = C_\nu \frac{k^2}{\varepsilon}
\end{aligned}$$

Here P – turbulent energy production an is given by (23), задається соотношением (), ε - rate of k dissipation, $Diffk_{urb}$ и $Diff\varepsilon_{urb}$ - turbulent diffusion of k and ε , depends on concrete model, ν_t – turbulent viscosity, C_1, C_2, C_ν - model constants.

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Анотація

Б. Головня

ПОВНА $k-\epsilon$ МОДЕЛЬ ТУРБУЛЕНТНОСТІ В ЦИЛІНДРИЧНИХ КООРДИНАТАХ

Ламінарні рівняння Нав'є-Стокса в циліндричних координатах описані в більшості книг з теоретичної гідродинаміки [1] - [3]. Рівняння Рейнольдса також легко знайти в літературі. Але повні рівняння $k-\epsilon$ моделі турбулентності в циліндричних координатах зустрічаються дуже рідко. Як правило, ці рівняння наводяться в роботах лише в спрощеному вигляді. У даній статті отримана не спрощена система рівнянь $k-\epsilon$ моделі турбулентності в циліндричних координатах

Ключові слова: *$k-\epsilon$ модель турбулентності, рівняння Рейнольдса в циліндричних координатах*

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